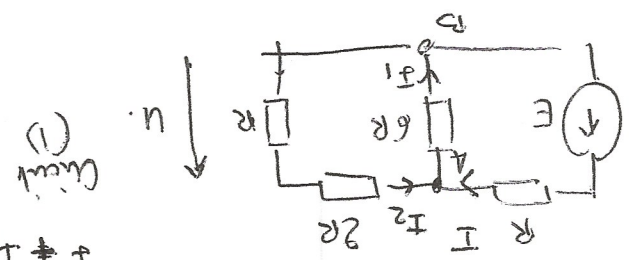


10/10/25

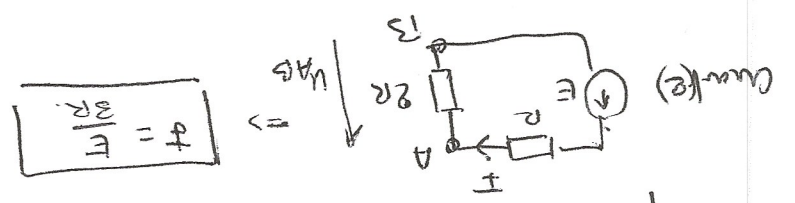
$I_1 = I_2 = I$ l'ingère l'équilibre de I_1 et I_2



10) a) 1^{er} méthode entre A et B, I_2 passant R et I_1 qui sort de la source

6R parcourue par I_1 et 3R parcourue par I_2 sort de la //

$$\frac{1}{R_{eq}} = \frac{1}{6R} + \frac{1}{3R} = \frac{3}{6R} \Rightarrow R_{eq} = 2R$$



$$I_1 = \frac{E}{3R}$$

entre A et B I_1 passant 2R et 3R parcourue par I_2 sort de la source

b) 2^{de} méthode du courant

$$\frac{I_1}{6R} + \frac{I_2}{3R} = \frac{1}{3} \Rightarrow$$

$$\frac{I_2}{3R} = \frac{2R}{3R} \Rightarrow I_2 = -\frac{2}{3}I$$

c) Vérification la 1^{re} méthode

$$I_1 + I_2 = I - \frac{2}{3}I = \frac{1}{3}I \text{ qui est bien égal à } I_1$$

d) $U = -RI_2 = -R(-\frac{2}{3}I) = +\frac{2}{3}RI = U$ la 1^{re} méthode est correcte

2^{de} méthode

a) $U = \frac{R}{3R} U_{AB}$ et ds le circuit (2) $U_{AB} = \frac{2R}{3R} \cdot E = \frac{2}{3}E$

d'après $U = \frac{2}{3}E$ On retrouve le résultat précédent.

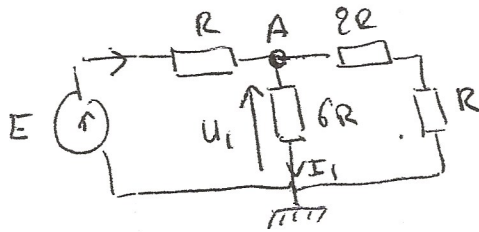
$$d'après $I_2 = -\frac{U}{R} = -\frac{\frac{2}{3}E}{R} = -\frac{2}{3}\frac{E}{R}$$$

$$I_1 = \frac{U_{AB}}{3R} = \frac{E}{3R}$$

$$I_1 + I_2 = \frac{E}{3R} - \frac{2E}{3R} = -\frac{E}{3R}$$

3^{ème} méthode

a- Pour alléger les calculs on fixe un point de masse et on le choisit



$$U_1 = V_A - V_M$$

||
0

$$V_A = \frac{\frac{E}{R} + 0 + 0}{\frac{1}{R} + \frac{1}{6R} + \frac{1}{3R}} = \frac{E}{\left(1 + \frac{1}{6} + \frac{1}{3}\right)} = \frac{6 \cdot E}{6 + 1 + 2} = \frac{6E}{9} = \boxed{\frac{2E}{3} = V_A} = U_1$$

$$\text{d'où } I_1 = \frac{U_1}{6R} = \frac{2E}{18R} = \boxed{\frac{E}{9R} = I_1} \quad \text{et} \quad \boxed{I_2 = -\frac{U_1}{3R} = -\frac{2E}{9R}}$$

AN

$$I = \frac{E}{3R} = \frac{3}{3 \cdot 1,5 \cdot 10^3} = \frac{2}{3} \cdot 10^{-3} \text{ A} \Rightarrow \boxed{I = 0,66 \text{ mA}}$$

$$I_1 = \frac{E}{9R} = \frac{I}{3} = \boxed{0,22 \text{ mA} = I_1}$$

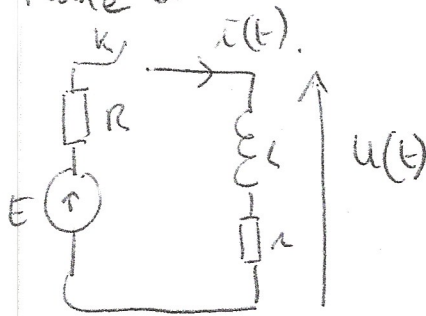
$$\boxed{I_2 = -0,44 \text{ mA}}$$

$$U = \frac{2}{9} \cdot E$$

$$= \frac{6}{9} = \frac{2}{3} = \boxed{0,66 \text{ V} = U}$$

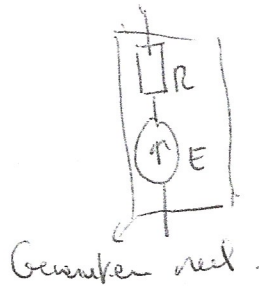
Ex 2

Partie 1.



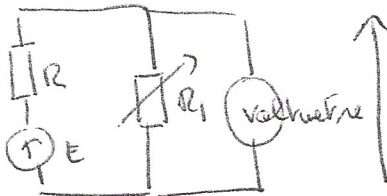
Question préliminaire: Détermination de la "force motrice"

1^{re} mesure. → On mesure $E =$ force à vide à l'aide d'un voltmètre de résistance interne $R_v \gg R$.



$$u_{mes} = \frac{R_v}{R + R_v} \cdot E \approx E$$

2^e mesure.



$$u_{mes} = \frac{R_1}{R + R_1} \cdot E$$

Pour $R_1 = R \Rightarrow u_{mes} = E/2$

On relève la valeur de R_1

pour en déduire celle de E .

1] K fermé. $E = (R + r)i + L \frac{di}{dt} \rightarrow \frac{di}{dt} + \frac{R+r}{L} i = \frac{E}{L}$

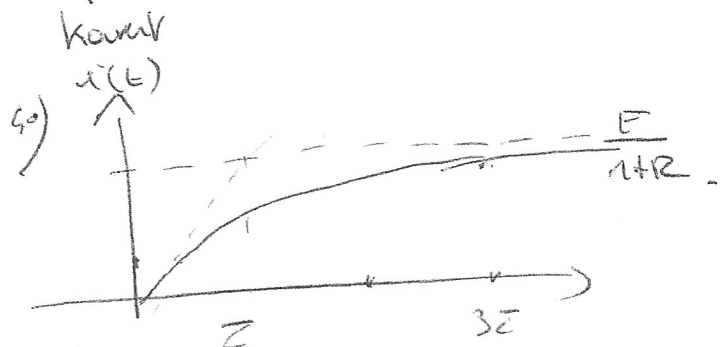
2] $\tau = \frac{L}{R+r}$ $i(0) \rightarrow i(\infty)$ régime permanent. et $i(\infty) = I_p = \frac{E}{R+r}$

AN $\tau = \frac{0.1}{60} = 1.6 \text{ ms}$

3] $i(t) = A e^{-\frac{t}{\tau}} + \frac{E}{R+r}$

Condition de i ds la bobine.
a. $i(0) = 0 \Rightarrow i(0^+) = 0$.

$$i(t) = \frac{E}{R+r} (1 - e^{-\frac{t}{\tau}})$$

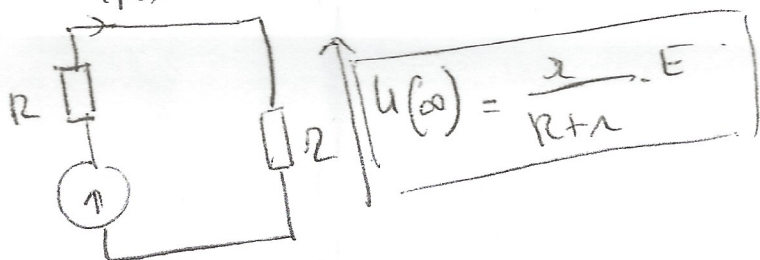


5.] da bei der nullen

ab $t=0^+$ $E = Ri + u$ $\bar{a}$ l'instant $t=0^+$ donne

$$u(0^+) = E$$

en regime permanent $m = \text{---}$



$$b]. u = Ri + L \frac{di}{dt} = \frac{R \cdot E}{R+L} \left(1 - e^{-\frac{t}{\tau}}\right) + L \cdot \frac{E}{R+L} \cdot \frac{1}{\tau} e^{-\frac{t}{\tau}}$$

$$= \frac{R}{R+L} \cdot E - \frac{R \cdot E}{R+L} e^{-\frac{t}{\tau}} + E e^{-\frac{t}{\tau}} = \frac{R}{R+L} \cdot E + E \left(1 - \frac{R}{R+L}\right) e^{-\frac{t}{\tau}}$$

$$u(t) = \frac{R}{R+L} \cdot E + \frac{L}{R+L} \cdot E e^{-\frac{t}{\tau}}$$

at $t=0$, e est nul

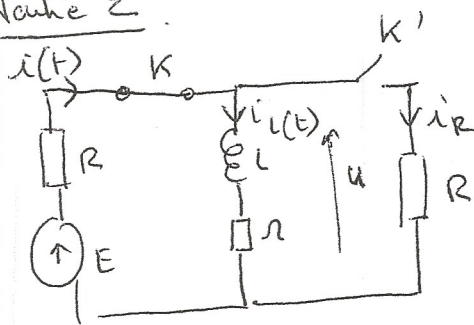
$$u(0^+) = \frac{R+L}{R+L} \cdot E = E$$

at $t \rightarrow \infty$

$$u(\infty) = \frac{R}{R+L} \cdot E$$

resultats du Sa
Ser hier un fus

Partie 2



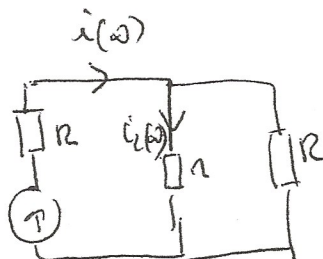
1- à $t=0$, on ferme K'

$t=0^-$, régime permanent établi $\Rightarrow$

$$i_L(0^-) = \frac{E}{R+r} \text{ et par continuité de } i$$

$$\text{ds } L \Rightarrow \boxed{i_L(0^+) = \frac{E}{R+r}}$$

e régime permanent



$$i(\infty) = \frac{E}{R + \frac{rR}{1+r}} = \frac{R+r \cdot E}{R^2 + 2rR}$$

Diviser de courant

$$\frac{i_L(\infty)}{i(\infty)} = \frac{\frac{rR}{1+r}}{R} = \frac{R}{1+R}$$

$$\Rightarrow \boxed{i_L(\infty) = \frac{R \cdot E}{R(R+2r)}}$$

2- équations diff à t

$$\begin{cases} E = Ri + u & (1) \\ i = i_L + i_R & (2) \\ u = L \frac{di_L}{dt} + r i_L = R i_R & (3) \end{cases}$$

$$\textcircled{1} \rightarrow \frac{E-u}{R} = i_L + \frac{u}{R} \rightarrow \frac{E}{R} = i_L + \frac{2u}{L} \Rightarrow \frac{E}{R} = i_L + \frac{2}{L} \left(L \frac{di_L}{dt} + r i_L \right)$$

$$\Rightarrow \frac{di_L}{dt} + \frac{R+2r}{2L} i_L = \frac{E}{2L}$$

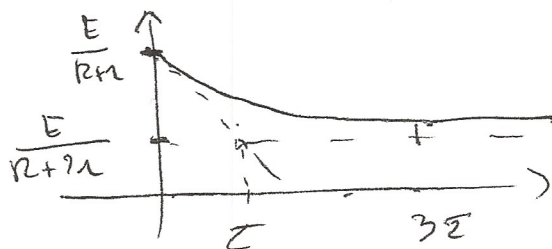
$$\boxed{\tau = \frac{2L}{R+2r}}$$

$$I_p = \frac{E}{R+2r} = i_L(\infty)!$$

$$i_L(t) = A e^{-\frac{t}{\tau}} + \frac{E}{R+2r}$$

$$i_L(0^+) = \frac{E}{R+r} = A + \frac{E}{R+2r} \rightarrow A = \frac{E}{R+r} - \frac{E}{R+2r} = E \left[\frac{R+2r - R - r}{(R+r)(R+2r)} \right]$$

$$\boxed{i_L(t) = \frac{E}{R+2r} \left[1 + \frac{r}{R+r} e^{-\frac{t}{\tau}} \right]}$$



exercice 3

10] $C(V_1 - V_2) = Q$ (Définition de C et la charge de l'armature au potentiel V_1)

$$\Rightarrow \boxed{C_0 = \frac{2\pi\epsilon_0 \cdot H}{\ln\left(\frac{a_2}{a_1}\right)}}$$

20] On pose $C_{milieu} = \frac{2\pi\epsilon_{milieu} \cdot R_{milieu}}{\ln\frac{a_2}{a_1}}$

la charge Q se répartit sur toute l'armature 1 et

$$Q = Q_{eau} + Q_{air} = \underbrace{\left[\frac{2\pi\epsilon_{eau} \cdot R}{\ln\frac{a_2}{a_1}} \right]}_{C_{eau}} + \underbrace{\left[\frac{2\pi\epsilon_{air} (H-R)}{\ln\frac{a_2}{a_1}} \right]}_{C_{air}} (V_1 - V_2)$$

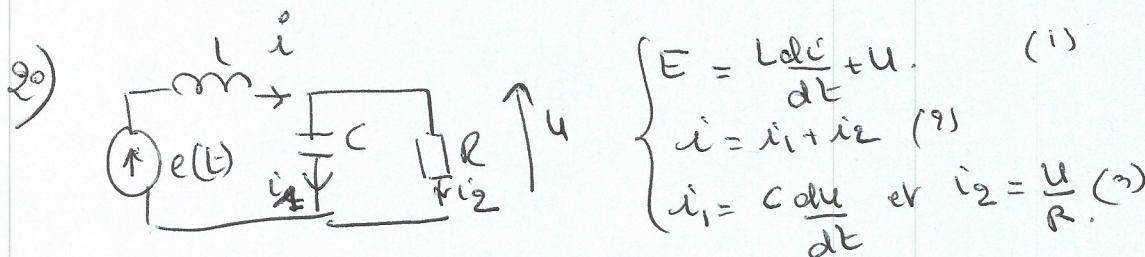
$$\Rightarrow \boxed{C = C_{eau} + C_{air}}$$

30] $C(R) = \frac{2\pi\epsilon_0}{\ln\frac{a_2}{a_1}} \left[\epsilon_r \cdot R + (H-R) \right] = \underbrace{\frac{2\pi\epsilon_0}{\ln\frac{a_2}{a_1}} \cdot H}_{C_0} \left[\frac{H-R}{H} + \epsilon_r \frac{R}{H} \right]$

$$C(R) = C_0 \cdot \left[\left(\frac{\epsilon_r - 1}{H} \right) \cdot R + 1 \right] = \boxed{C_0 (AR + B) = C(R)}$$

avec $\boxed{A = \frac{\epsilon_r - 1}{H} \quad B = 1}$

Cas $R=0$ $C = C_0$ (Absence d'eau)
 $R=H$ $AH + B = \epsilon_r - 1 + 1 \Rightarrow C = C_{eau} = \frac{2\pi\epsilon_0 \epsilon_r R}{\ln\frac{a_2}{a_1}}$ (Absence d'air)



D'après (2) et (3) $i = C \frac{du}{dt} + \frac{u}{R}$ et (1) $\Rightarrow E = L \left[C \frac{d^2u}{dt^2} + \frac{1}{R} \frac{du}{dt} \right] + u$

$$E = LC \frac{d^2 u}{dt^2} + \frac{L}{R} \frac{du}{dt} + u \Rightarrow \frac{d^2 u}{dt^2} + \frac{1}{RC} \frac{du}{dt} + \frac{1}{LC} u = \frac{E}{LC}$$

On pose $\omega_0^2 = \frac{1}{LC} \Rightarrow \boxed{\omega_0 = \frac{1}{\sqrt{LC}}}$ $\frac{\omega_0}{Q} = \frac{1}{RC} \rightarrow Q = RC\omega_0 \Rightarrow$

$Q = RC \cdot \frac{1}{\sqrt{LC}} \rightarrow \boxed{Q = R \sqrt{\frac{C}{L}}}$ $A = E$

* ω_0 = pulsation propre du circuit = pulsation des oscillations sans amortissement $\Rightarrow Q \rightarrow \infty$ ($\lambda = 0$) et $R \rightarrow \infty$. (marque avant).

* Q facteur de qualité $\Rightarrow$ permet de distinguer les régimes $Q = \frac{1}{2\lambda} \rightarrow \lambda = \frac{1}{2Q} \Rightarrow \boxed{\lambda = \frac{1}{2R} \sqrt{\frac{L}{C}}}$

λ facteur d'amortissement.

$$\begin{array}{ccc} [\omega^2][u] = \left[\frac{d^2 u}{dt^2} \right] = \frac{[E]}{[Q]} \left[\frac{du}{dt} \right] & \Rightarrow & [\omega^2] = T^{-2} \Rightarrow [u] = T^{-1} \\ \downarrow & & \downarrow \\ \text{Temps} & & \text{Temps}^2 \end{array} \quad \begin{array}{c} \text{Temps} \\ \text{Temps}^2 \end{array}$$

$\Rightarrow [\omega^2] = T^{-2} \Rightarrow [u] = T^{-1}$
et $\left[\frac{\omega_0}{Q} \right] = T^{-1} \Rightarrow [Q] = 1$

2°] Résolution de l'équation

$$u(t) = S_{ho}(t) + S_p = \underbrace{S_{ho}(t)}_{\text{Solution de l'équation homogène}} + E$$

Résolution de l'eq. homogène:

$$\frac{d^2 u}{dt^2} + \frac{\omega_0}{Q} \frac{du}{dt} + \omega_0^2 u = 0$$

et solution $\Rightarrow r^2 + \frac{\omega_0}{Q} r + \omega_0^2 = 0$

$$\Delta = \frac{\omega_0^2}{Q^2} - 4\omega_0^2 = \omega_0^2 \left[\frac{1}{Q^2} - 4 \right]$$

le cas $\Delta > 0$ $\left[\frac{1}{Q^2} - 4 \right] > 0 \rightarrow Q < \frac{1}{2}$ $r_{1,2} = -\frac{\omega_0}{2Q} \pm \frac{1}{2} \cdot \sqrt{\omega_0^2 \left(\frac{1}{Q^2} - 4 \right)}$

Regime aperiodique
 $\Rightarrow r_{1,2} = -\frac{\omega_0}{2Q} \pm \frac{\omega_0}{2Q} \sqrt{1 - 4Q^2}$

$$\boxed{u(t) = A e^{r_1 t} + B e^{r_2 t}}$$

$$\Delta = 0 \quad \varphi = \frac{1}{2} \text{ regime critique} \quad r_1 = r_2 = -\omega_0$$

$$u(t) = (A + Bt) e^{-\omega_0 t}$$

$$\underline{\Delta < 0} \quad \varphi > \frac{1}{2} \quad r_{1,2} = -\frac{\omega_0}{2\varphi} \pm i\omega_0 \sqrt{1 - \frac{1}{4\varphi^2}}$$

$$\boxed{u(t) = A e^{-\frac{\omega_0}{2\varphi} t} \cos(\omega t + \varphi) \text{ avec } \omega = \omega_0 \sqrt{1 - \frac{1}{4\varphi^2}}}$$

3°] Déterminer de $u(0^+)$ et $\frac{du}{dt}(0^+)$

à $t=0^-$ le condensateur est déchargé $\Rightarrow u(0^-) = 0$.

$$u(0^+) = u(0^-) = 0 \Rightarrow \boxed{u(0^+) = 0}$$

↳ Continuité de u aux bornes de C .

$$i_1 = C \frac{du}{dt} \Rightarrow \frac{du}{dt}(0^+) = \frac{i_1(0^+)}{C}$$

Capacitance de $i_1(0^+)$ à $t=0^-$ $i(0^-) = 0$ (Aucun courant ne circule).

$$i(0^+) = i(0^-) = 0$$

↳ Continuité de i traversant une bobine.

$i(0^+) = 0$

$i_2(0^+) = 0 = \frac{u(0^+)}{R}$

$\Rightarrow i_1(0^+) = 0 \Rightarrow \boxed{\frac{du}{dt}(0^+) = 0}$

$$* \quad u(t) = A e^{-\frac{\omega_0}{2\varphi} t} \cos(\omega t + \varphi) + E \quad \omega = \omega_0 \sqrt{1 - \frac{1}{4\varphi^2}} \quad \frac{\omega_0}{\varphi} = \frac{1}{RC}$$

$$u(0^+) = A \cos \varphi + E = 0 \rightarrow A = -\frac{E}{\cos \varphi}$$

$$* \quad \frac{du}{dt} = A e^{-\frac{\omega_0}{2\varphi} t} \left[-\frac{\omega_0}{2\varphi} \cos(\omega t + \varphi) - \omega_0 \sqrt{1 - \frac{1}{4\varphi^2}} \sin(\omega t + \varphi) \right]$$

$$\frac{du}{dt}(0^+) = A e^{-\frac{\omega_0}{2\varphi} t} \left[-\frac{\omega_0}{2\varphi} \cos \varphi - \omega_0 \sqrt{1 - \frac{1}{4\varphi^2}} \sin \varphi \right] = 0$$

$$\frac{\cos \varphi}{2\varphi} + \sqrt{1 - \frac{1}{4\varphi^2}} \sin \varphi = 0$$

$$\tan \varphi = -\frac{1}{2\varphi \sqrt{1 - \frac{1}{4\varphi^2}}} \quad \text{et} \quad \frac{1}{\cos^2 \varphi} = 1 + \tan^2 \varphi = 1 + \frac{\frac{1}{4\varphi^2}}{1 - \frac{1}{4\varphi^2}} = \frac{1}{1 - \frac{1}{4\varphi^2}}$$

$$\left\{ \begin{array}{l} \cos \varphi = \sqrt{1 - \frac{1}{4\varphi^2}} \\ \varphi = \arccos \left[\frac{-\frac{1}{2\varphi}}{\sqrt{1 - \frac{1}{4\varphi^2}}} \right] \end{array} \right. \quad -\frac{\pi}{2} < \varphi < +\frac{\pi}{2}$$

$$\Rightarrow A = -\frac{E}{\sqrt{1 - \frac{1}{4\varphi^2}}} \quad u(t) = E \left(1 - \frac{e^{-\frac{\omega_0 t}{2\varphi}}}{\sqrt{1 - \frac{1}{4\varphi^2}}} \cos(\omega t + \varphi) \right)$$

$$4) \quad \delta = \ln \frac{u(t) - u_{\infty}}{u(t+\tau) - u_{\infty}} = \ln \frac{A e^{-\frac{\omega_0 t}{2\varphi}} \cos(\omega t + \varphi)}{A e^{-\frac{\omega_0 (t+\tau)}{2\varphi}} \cos(\omega (t+\tau) + \varphi)} = \frac{\omega_0 \tau}{2\varphi}$$

$$\delta = \frac{\omega_0 \tau}{2\varphi} = \frac{2\pi}{2\varphi \cos \varphi \sqrt{1 - \frac{1}{4\varphi^2}}} \Rightarrow \left[\delta = \frac{1}{2\varphi \sqrt{1 - \frac{1}{4\varphi^2}}} \right] \quad \text{ou} \quad \left[\delta = \frac{2\pi}{\sqrt{4\varphi^2 - 1}} \right]$$

$$\delta = \ln \frac{3}{1} = 1,1 = \delta \quad \varphi = \frac{1}{2} \sqrt{\frac{4\pi^2}{\delta^2} + 1} \Rightarrow \boxed{\varphi = 2,9}$$

$$T = 1 \mu s = \frac{T_0}{\sqrt{1 - \frac{1}{4\varphi^2}}} \Rightarrow T_0 = T \sqrt{1 - \frac{1}{4\varphi^2}} = 10^{-6} \sqrt{1 - \frac{1}{4 \cdot (2,9)^2}} \\ \boxed{T_0 = 0,97 \mu s}$$

$$5) \quad \left[h = \left(\frac{C}{C_0} - 1 \right) \frac{H}{\epsilon_1 - 1} \right] \quad T_0 = 2\pi \sqrt{LC} \Rightarrow C = \frac{T_0^2}{4\pi^2 L} \quad \text{AN} \quad \boxed{C = 4,2 \cdot 10^{-11} \text{ F}}$$

$$C_0 = \frac{2\pi \epsilon_0 H}{\ln\left(\frac{a_2}{a_1}\right)} \quad \text{AN} \rightarrow \boxed{C_0 = 0,5 \cdot 10^{-11} \text{ F}} \quad h = \left(\frac{4,2}{0,5} - 1 \right) \cdot \frac{10}{\epsilon_0} = \boxed{0,9 \text{ cm} = h}$$

$$\boxed{V = \pi R (a_2^2 - a_1^2) = 2,5 \cdot 10^{-5} \text{ m}^3}$$